Dechao Zheng

Chongqing University

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Chongqing University and Vanderbilt University

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Spectrum

The spectrum $\sigma(T)$ of a bounded linear operator T acting on a Hilbert space H is the set of complex numbers λ such that $\lambda I - T$ does not have an inverse that is a bounded linear operator.

If $H = C^n$, T can be viewed as a matrix

$$\begin{bmatrix} a_{11} & a_{12} & \cdots & a_{1n} \\ a_{21} & a_{22} & \cdots & a_{2n} \\ \vdots & \vdots & \cdots & \vdots \\ a_{n1} & a_{n2} & \cdots & a_{nn} \end{bmatrix}$$

and so its spectrum consists of eigenvalues of the matrix. But if H is an infinite dimensional Hilbert space, the spectrum of its bounded operator T may have more numbers than its eigenvalues $\sigma_p(T)$.

ESSENTIAL SPECTRUM

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$$\sigma_e(T) = \{\lambda \in C : [\lambda I - T] \text{ is not invertible in } C(H).\}$$

Calkin algebra C(H) = B(H)/K(H)

B(H): the algebra of bounded linear operators on H.

K(H): the ideal of compact operators on H.

Fredholm Index and Spectral Picture

If λ is not in $\sigma_e(T)$, $T-\lambda I$ is Freholm. The Fredholm index is defined by

$$ind(T - \lambda I) = \dim Ker(T - \lambda I) - \dim Ker(T - \lambda I)^*.$$

FREDHOLM INDEX AND SPECTRAL PICTURE

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Carl Pearcy, Some Recent developments in Operator theory, CBMS 36, 1975.

THEOREM

Let Ω be a connected component of $C \setminus \sigma_e(T)$ such that $\operatorname{ind}(T - \lambda I) = 0$ for each $\lambda \in \Omega$. Then one of the following holds:

- (a) $\Omega \cap \sigma(T)$ is empty.
- (b) $\Omega \subset \sigma(T)$.
- (c) $\Omega \cap \sigma(T)$ is a countable set of isolated eigenvalues of T, each having finite multiplicity.

Furthermore the intersection of $\sigma(T)$ with the unbounded component of $C \setminus \sigma_e(T)$ is a countable set of isolated eigenvalues of T, each of which has finite multiplicity.

TOEPLITZ OPERATORS ON THE HARDY SPACE

A Toeplitz operator on the Hardy space is the compression of a multiplication operator on the circle to the Hardy space

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Let ∂D be the circle, with the standard Lebesgue measure, and $L^2(\partial D)$ be the Hilbert space of square-integrable functions. A bounded measurable function ϕ on ∂D defines a multiplication operator M_ϕ on $L^2(\partial D)$. Let P be the projection from $L^2(\partial D)$ onto the Hardy space H^2 . The Toeplitz operator with symbol ϕ is defined by

$$T_{\phi} = PM_{\phi}|_{H^2}$$

Toeplitz matrix

A bounded operator on H^2 is Toeplitz if and only if its matrix representation, in the basis $\{z^n\}_0^\infty$, has constant diagonals:

$$\begin{bmatrix} a_0 & a_{-1} & a_{-2} & a_{-3} & \cdots & \cdots \\ a_1 & a_0 & a_{-1} & a_{-2} & \cdots & \cdots \\ a_2 & a_1 & a_0 & a_{-1} & \cdots & \cdots \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \end{bmatrix}$$

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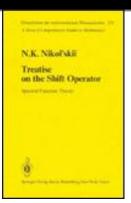
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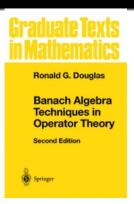
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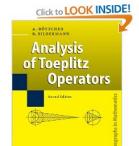
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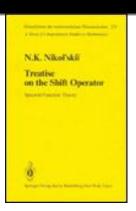


Spring









BERGMAN SPACE

Let dA denote Lebesgue area measure on the unit disk \mathbb{D} , normalized so that the measure of \mathbb{D} equals 1. The Bergman space L_a^2 is the Hilbert space consisting of the analytic functions on \mathbb{D} that are also in $L^2(\mathbb{D}, dA)$:

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$$f(z) = \sum_{n=0}^{\infty} a_n z^n,$$

and

$$\sum_{n=0}^{\infty} \frac{|a_n|^2}{n+1} < \infty.$$

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$$f(z) = \sum_{n=0}^{\infty} a_n z^n,$$

and

$$\sum_{n=0}^{\infty} \frac{|a_n|^2}{n+1} < \infty.$$

Let $e_n = \sqrt{n+1}z^n$. Then $\{e_n\}_0^\infty$ form an orthonormal basis of the Bergman space L_a^2 .

Toeplitz Operators

For $\phi \in L^{\infty}(\mathbb{D}, dA)$ where dA is normalized area measure on \mathbb{D} , the Toeplitz operator T_{ϕ} with symbol ϕ is the operator on L^2_a defined by

$$T_{\phi}f = P(\phi f);$$

here P is the orthogonal projection from $L^2(\mathbb{D},dA)$ onto L^2_a . Note that if $\phi \in H^\infty$ (the set of bounded analytic functions on $\partial \mathbb{D}$), then T_ϕ is just the operator of multiplication by ϕ on L^2_a .

Operator Theory in Function Spaces

Kehe Zhu



MATRIX REPRESENTATION

Let
$$e_n = \sqrt{n+1}z^n$$
 and $\phi(z) = \sum_{j=-\infty}^{-1} a_j \bar{z}^{|j|} + \sum_{j=0}^{\infty} a_j z^j$.

$$\langle T_{\phi} e_i, e_j \rangle = \sqrt{i+1} \sqrt{j+1} a_{j-i} \langle z^j, z^j \rangle = a_{j-i} \sqrt{\frac{i+1}{j+1}}.$$

On the basis $\{e_n\}$, the matrix representation of the Toeplitz operator \mathcal{T}_ϕ on the Bergman space is given by

$$\begin{bmatrix} a_0 & \sqrt{\frac{2}{1}} a_{-1} & \sqrt{\frac{3}{1}} a_{-2} & \sqrt{\frac{4}{1}} a_{-3} & \cdots & \cdots \\ \sqrt{\frac{1}{2}} a_1 & a_0 & \sqrt{\frac{3}{2}} a_{-1} & \sqrt{\frac{4}{2}} a_{-2} & \cdots & \cdots \\ \sqrt{\frac{1}{3}} a_2 & \sqrt{\frac{2}{3}} a_1 & a_0 & \sqrt{\frac{4}{3}} a_{-1} & \cdots & \cdots \\ \vdots & \ddots & \ddots & \ddots & \vdots \end{bmatrix}$$

SOME ALGEBRAIC PROPERTIES

(a)
$$T_{\alpha\phi+\beta\psi} = \alpha T_{\phi} + \beta T_{\psi}$$
.

(b) If ϕ is in H^{∞} , then

$$T_{\psi}T_{\phi}=T_{\psi\phi}.$$

(c) If $\overline{\psi}$ is in H^{∞} , then

$$T_{\psi}T_{\phi}=T_{\psi\phi}.$$

- (d) $T_{\phi}^* = T_{\overline{\phi}}$.
- (e) If $\phi \geq 0$, then $T_{\phi} \geq 0$.

Fredholm index for Toeplitz operator

If ϕ is continuous on the unit circle ∂D and does not vanish on ∂D , then T_ϕ is Fredholm and

$$ind(T_{\phi}) = n(\phi(\partial \mathbb{D}), 0).$$

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For a closed curve γ in the complex plane $\mathbb C$ and $a\in\mathbb C\backslash\gamma$, define the winding number of the curve γ with respect to a to be

$$n(\gamma, a) = \frac{1}{2\pi i} \int_{\gamma} \frac{dz}{z - a}.$$

BERGMAN SHIFT

On the basis $\{e_n = \sqrt{n+1}z^n\}$, the Toeplitz operator T_z with symbol z is a weighted shift operator, called the Bergman shift:

$$T_z e_n = \sqrt{\frac{n+1}{n+2}} e_{n+1},$$

and hence $T_{\bar{z}}$ is a backward weighted shift:

$$T_{\bar{z}}e_n = \begin{cases} 0 & n = 0\\ \sqrt{\frac{n}{n+1}}e_{n-1}. & n > 0 \end{cases}$$
 (1)

The matrix representation of the Toeplitz operators $T_{1-|z|^2} = I - T_z^* T_z$ is given by

$$\begin{bmatrix} \frac{1}{2} & 0 & 0 & 0 & \cdots & \cdots \\ 0 & \frac{1}{3} & 0 & 0 & \cdots & \cdots \\ 0 & 0 & \frac{1}{4} & 0 & \cdots & \cdots \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \end{bmatrix}$$



THEOREM (Coburn Theorem)

If $T_{\phi} \neq 0$ on the Hardy space, either $\ker T_{\phi} = \{0\}$ or $\ker T_{\phi}^* = \{0\}$.

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Does Coburn theorem hold on the Bergman space?

No! On the Bergman space, both $kerT_{1-|z|^2-\frac{1}{2}}$ and $kerT_{1-|z|^2-\frac{1}{2}}^*$ contain the function 1.

But $1-|z|^2-\frac{1}{2}$ is not harmonic in the unit disk and $T_{1-|z|^2}$ is compact!

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WIDOM THEOREM AND DOUGLAS THEOREM

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THEOREM (Douglas Theorem)

The essential spectrum $\sigma_e(T_\phi)$ of a Toeplitz operator on the Hardy space is also connected.



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COMPACT TOEPLITZ OPERATORS

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 $T_{1-|z|^2}$ is compact with the spectrum $\{\frac{1}{2},\frac{1}{3},\cdots\}\cup\{0\}$. Hence $\sigma(T_{1-|z|^2})$ is disconnected.

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 $T_{1-|z|^2}$ is compact with the spectrum $\{\frac{1}{2},\frac{1}{3},\cdots\}\cup\{0\}$. Hence $\sigma(T_{1-|z|^2})$ is disconnected. But $1-|z|^2$ is not harmonic on the unit disk.

Compact Toeplitz operators on the Hardy space and the Bergman space

THEOREM

On the Hardy space, T_{ϕ} is compact if and only if $\phi = 0$.

THEOREM (Axler-Zheng)

For $\phi \in L^{\infty}(D)$, T_{ϕ} is compact on the Hardy space if and only if

$$\lim_{|z|\to 1} \int_D \phi(w) \frac{(1-|z|^2)^2}{|1-\bar{z}w|^4} dA(w) = 0.$$

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If ϕ is harmonic on the unit disk, then

$$\lim_{|z|\to 1} \int_D \phi(w) \frac{(1-|z|^2)^2}{|1-\bar{z}w|^4} dA(w) = 0$$

implies that $\phi = 0$ on ∂D and hence $\phi = 0$ on the unit disk.

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$$\lim_{|z| \to 1} \int_{D} \phi(w) \frac{(1 - |z|^{2})^{2}}{|1 - \overline{z}w|^{4}} dA(w) = 0$$

implies that $\phi=0$ on ∂D and hence $\phi=0$ on the unit disk. There is no nontrivial compact Toeplitz operator with bounded harmonic symbol on the Bergman space.

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Is the spectrum $\sigma(T_{\phi})$ of a Toeplitz operator on the Bergman space connected if ϕ is bounded and harmonic on the unit disk?

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Sundberg's conjecture: Yes! (Problem 7.10 in V.P. Havin and N.K. Nikolski (Eds), Linear and Complex Analysis Problem Book 3, Lecture notes in Mathematics 1573, 1994).

7.10

TOEPLITZ OPERATORS ON THE BERGMAN SPACE

C. SUNDBERG

Let A^2 denote the Bergman space of analytic functions in $L^2(\mathbb{D})$, and let P be the orthogonal projection of $L^2(\mathbb{D})$ onto A^2 . For $\varphi \in L^\infty(\mathbb{D})$ we define the Toeplitz operator with symbol φ by $T_{\alpha} = P(\varphi f)$. In general the behaviour of these operators may be quite different from that of the Toeplitz operators on the Hardy space H^2 . However it is shown in [1] that Toeplitz operators on A2 with harmonic symbols behave quite similarly to those on H^2 , and one can prove analogues for this class of many results about Toeplitz operators on H^2 .

An important result about Toeplitz operators on H² is Widom's Theorem, which states that the spectrum of such an operator is connected ([2]). This suggests our

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This was an open question in (G. McDonald and C. Sundberg, *Indiana Univ. Math. J.* 28 (1979)).

Let ϕ be in $H^{\infty}(D)$.

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If λ is not in the closure of $\phi(D)$, then $\frac{1}{\phi-\lambda}$ is in $H^\infty(D)$ and

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If $\lambda = \phi(a)$ for some $a \in D$, then

$$T_{\phi-\lambda}^* k_a = 0.$$

Hence

(a) If ϕ is analytic on the unit disk, then

$$\sigma(T_{\phi}) = clos\phi(D).$$



(b) If ϕ is real and harmonic on the unit disk, then

$$\sigma(T_{\phi}) = [\inf \phi, \sup \phi].$$

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(c) If ϕ is harmonic and has piecewise continuous boundary values, then $\sigma_e(T_\phi)$ consists of the path formed boundary values of ϕ by joining the one-sided limits at discontinuities by straight line segments and hence $\sigma_e(T_\phi)$ is connected.

(b) If ϕ is real and harmonic on the unit disk, then

$$\sigma(T_{\phi}) = [\inf \phi, \sup \phi].$$

- (c) If ϕ is harmonic and has piecewise continuous boundary values, then $\sigma_e(T_\phi)$ consists of the path formed boundary values of ϕ by joining the one-sided limits at discontinuities by straight line segments and hence $\sigma_e(T_\phi)$ is connected.
- (b) and (c) are contained in (G. McDonald and C. Sundberg, *Indiana Univ. Math. J.* 28 (1979)).

We hope to construct ϕ having the following properties:

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- (b) h(z) is continuous on the closure of the unit disk.
- (c) $\sigma_e(T_h) = h(\partial D)$.
- (d) 0 is an isolated eigenvalue of T_h .

Eigenvectors of T_h for the eigenvalue 0

Let f be an eigenvector for T_h for the eigenvalue 0. Then

$$0 = T_h f(z)$$

$$= T_{\overline{z}} f(z) + T_{\phi(z)} f(z)$$

$$= \frac{1}{z^2} \int_0^z w f'(w) dw + \phi(z) f(z).$$

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LEMMA

For f in the Bergman space L_a^2 ,

$$T_{\bar{z}}f(z) = \frac{1}{z^2} \int_0^z w f'(w) dw.$$
 (2)

$$\frac{1}{z^2} \int_0^z w f'(w) dw + \phi(z) f(z) = 0$$

This is equivalent to the following first order differential equation

$$(1 + z\phi(z))f'(z) = -(2\phi(z) + z\phi'(z))f(z).$$
 (3)

For a fixed 0 < r < 1, we want

(a) a rational function $\eta(z)$ with poles outside the closure of the unit disk such that

$$2\phi(z) + z\phi'(z) = (z - r)\eta(z);$$

$$\frac{1}{z^2}\int_0^z wf'(w)dw + \phi(z)f(z) = 0$$

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(b) $1+z\phi(z)$ has a simple zero at z=r and no other zeros in $\overline{\mathbb{D}}$. Write

$$\psi(z) = \frac{1 + z\phi(z)}{z - r}.$$

Then ψ is a rational function with poles outside of the closure of the unit disk.

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(3) becomes

$$\frac{f'(z)}{f(z)} = -\frac{\eta(z)}{\psi(z)}.$$



$$(1 + z\phi(z))f'(z) = -(2\phi(z) + z\phi'(z))f(z)$$

A solution of the above equation in the Bergman space \mathcal{L}_a^2 is given by

$$f(z) = exp[-\int_0^z \frac{\eta(w)}{\psi(w)} dw].$$

Thus f is an eigenvector of T_h for the eigenvalue equal to 0.

$$(1 + z\phi(z))f'(z) = -(2\phi(z) + z\phi'(z))f(z)$$

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Thus f is an eigenvector of T_h for the eigenvalue equal to 0. Since $\sigma_e(T_h) = h(\partial D)$ and

$$ind(T_h) = n(h(\partial \mathbb{D}), 0),$$

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$$dim\ kerT_h = dim\ kerT_h^*$$
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LEMMA

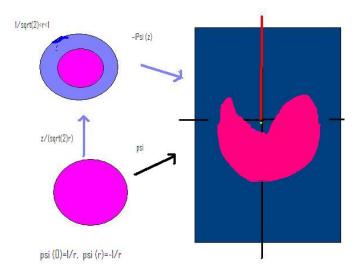
For each 0 < r < 1, there exists a rational function $\phi(z)$ with poles outside $\overline{\mathbb{D}}$ such that

- (a) $2\phi(r) + r\phi'(r) = 0$.
- (b) $1+z\phi(z)$ has a simple zero at z=r and no other zeros in $\overline{\mathbb{D}}$.
- (c) The winding number

$$n(h(\partial \mathbb{D}), 0) = 0$$

where $h = \overline{z} + \phi(z)$.

Sketch of Proof



DISCONNECTED SPECTRUM

THEOREM

Let $h(z) = \bar{z} + \phi(z)$ Then 0 is an eigenvalue of T_h and is an isolated point of $\sigma(T_h)$. Hence $\sigma(T_h)$ is disconnected.

(1) Since h is continuous on the closure of the unit disk, then

$$\sigma_e(T_h) = h(\partial \mathbb{D}).$$

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(2) $0 \in \sigma_p(T_h) \cap \Omega$ where

$$\Omega = \{\lambda \notin \sigma_e(T_h) : ind(T_h - \lambda I) = 0\}$$
$$= \{\lambda \notin h(\partial \mathbb{D}) : n(h(\partial \mathbb{D}), \lambda) = 0\}.$$

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Sketch of the proof Want: $\Omega \cap \sigma_p(T_h)$ is countable.



CAUCHY'S ARGUMENT PRINCIPLE

Cauchy's argument principle

if f(z) is a meromorphic function inside and on some closed contour C, and f has no zeros or poles on C, then

$$\oint_C \frac{f'(z)}{f(z)} dz = 2\pi i (N - P)$$

$$\lambda \in \Omega \cap \sigma_p(T_h), \ n(h(\partial \mathbb{D}), \lambda) = 0$$

Since $\frac{1}{z} + \phi(z)$ has a simple pole at z = 0 and no other poles in the unit disk \mathbb{D} , the argument principle tells us that if λ is in $\Omega \cap \sigma_p(T_h)$, there is a unique point z_λ in \mathbb{D} such that

$$\frac{1}{z_{\lambda}} + \phi(z_{\lambda}) = \lambda.$$

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As λ is an eigenvalue of T_h , there is a nonzero function g in the Bergman space L^2_a such that

$$\lambda g = T_h g(z)$$

$$= T_{\bar{z}} g(z) + T_{\phi(z)} g(z)$$

$$= \frac{1}{z^2} \int_0^z w g'(w) dw + \phi(z) g(z).$$

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We solve the above equation to obtain

$$\frac{g'(z)}{g(z)} = -\frac{2(\phi(z) - \lambda) + z\phi'(z)}{1 + z(\phi(z) - \lambda)}.$$

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$$-\frac{2(\phi(z_{\lambda})-\lambda)+z_{\lambda}\phi'(z_{\lambda})}{\phi(z_{\lambda})-\lambda+z_{\lambda}\phi'(z_{\lambda})}=-1-\frac{1}{1-z_{\lambda}^2\phi'(z_{\lambda})}.$$

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The regularity of g(z) at $z=z_{\lambda}$ forces this residue to be in $\mathbb{N}=\{0,1,2,3,\cdots,\}$ which leads to

$$z_{\lambda}^2\phi'(z_{\lambda})=1+\frac{1}{n+1}$$

for some $n \in \mathbb{N}$.

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for some $n \in \mathbb{N}$. This restricts the set

$$\Omega \cap \sigma_p(T_h) \subset \{\lambda : \lambda = \frac{1}{z_\lambda} + \phi(z_\lambda), z_\lambda^2 \phi'(z_\lambda) = 1 + \frac{1}{n+1}, \text{ for some } n \in \mathbb{N}\}$$

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to be a countable set. Thus 0 is an isolated point in $\sigma(T)$. Hence we conclude that the spectrum $\sigma(T_h)$ is disconnected.

For $z \in \mathbb{D}$, let ϕ_z be the analytic map of \mathbb{D} onto \mathbb{D} defined by

$$\phi_z(w) = \frac{z - w}{1 - \bar{z}w}.$$
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For S a bounded operator on L_a^2 , define S_z to be the bounded operator on L_a^2 given by conjugation with U_z :

$$S_z = U_z S U_z$$
.



MAXIMAL IDEAL SPACE OF H^{∞}

Let $\mathcal M$ be the maximal ideal space of H^∞ , i.e., the set of complex homomorphisms of H^∞ with w^* -topology. Then $\mathcal M$ is a compact Hausdorff space.

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Let \mathcal{M} be the maximal ideal space of H^{∞} , i.e., the set of complex homomorphisms of H^{∞} with w^* -topology. Then \mathcal{M} is a compact Hausdorff space. If z is a point in the unit disk \mathbb{D} , then point evaluation at z is a multiplicative linear functional on \mathcal{M} . Thus we can think of z as an element of \mathcal{M} and the unit disk \mathbb{D} as a subset of \mathcal{M} . Carleson's corona theorem states that \mathbb{D} is dense in \mathcal{M} .

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$$\lim_{z \to m} \alpha_z = \beta$$

means (as you should expect) that for each open set X in E containing β , there is an open set Y in \mathcal{M} containing m such that $\alpha_z \in X$ for all $z \in Y \cap \mathbb{D}$. Note that with this notation z is always assumed to lie in \mathbb{D} .

HOFFMAN MAP

For $m \in \mathcal{M}$, let $\phi_m \colon \mathbb{D} \to \mathcal{M}$ denote the *Hoffman map*. This is defined by setting

$$\phi_m(w) = \lim_{z \to m} \phi_z(w)$$

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for $w \in \mathbb{D}$; here we are taking a limit in \mathcal{M} . The existence of this limit, as well as many other deep properties of ϕ_m , was proved by Hoffman (*Ann. Math.*, **103** (1967)).

Localization S_m of S in Toeplitz algebra

The Toeplitz algebra $\mathcal T$ is the C^* -subalgebra of $\mathcal B(L^2_a)$ generated by $\{T_g:g\in H^\infty\}.$

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LEMMA

If $S \in \mathcal{T}$, the Toeplitz algebra and $m \in \mathcal{M}$, then there exists $S_m \in \mathcal{T}$ such that

$$\lim_{z\to m}\|S_zf-S_mf\|=0$$

for every f in L_a^2 . If $S = T_{u_1} \dots T_{u_n}$, where $u_1, \dots, u_n \in \mathcal{U}$, then $S_m = T_{u_1 \circ \phi_m} \dots T_{u_n \circ \phi_m}$.

ESSENTIAL SPECTRUM

Using a similar argument as one in the proof of Theorem 10.3 in (D. Suarez, *Indiana Univ. Math. J.*, **56** (2007)), we have the following theorem.

THEOREM

If $S \in \mathcal{T}$, the Toeplitz algebra, then

$$\mathbb{C}ackslash\sigma_e(S)=\{\lambda\in\mathbb{C}:\lambda
otin\mathbb{D}\ \sigma(S_m)\ \ ext{and}$$

$$\sup_{m\in\mathcal{M}\setminus\mathbb{D}}\|(S_m-\lambda I)^{-1}\|<\infty\}.$$

THIN BLASCHKE PRODUCT

To a sequence $\{z_n\}_n$ in $\mathbb D$ with $\sum_{n=1}^\infty (1-|z_n|)<\infty$, there corresponds a Blaschke product

$$b(z) = \prod_{n=1}^{\infty} \frac{-\overline{z}_n}{|z_n|} \frac{z - z_n}{1 - \overline{z}_n z}, \quad z \in \mathbb{D}.$$

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A sequence $\{z_n\}_n$ and its associated Blaschke product are called thin if

$$\lim_{n\to\infty}\prod_{k\neq n}\left|\frac{z_n-z_k}{1-\overline{z}_kz_n}\right|=1.$$

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Hedenmalm (*Proc. Amer. Math. Soc.*, 99 (1987)) showed that for each m in $\mathcal{M}\backslash\mathbb{D}$, either

$$b \circ \phi_m(z) = \lambda_m$$
 or $b \circ \phi_m(z) \in Aut(\mathbb{D})$

for some unimodular constant λ_m . The latter case actually occurs if m is in the Gleason part of some point in the closure of zeros of b in \mathbb{D} .

THEOREM

Let F be a continuous function on the closure $\overline{\mathbb{D}}$ of the unit disk, b be an infinite thin Blaschke product and $F_b = F \circ b$. Then

$$\sigma_e(T_{F_b}) = \sigma(T_F).$$

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Proof Let $S = T_{F_b}$. For each m in $\mathcal{M} \setminus \mathbb{D}$,

$$S_m = T_{F \circ b \circ \phi_m}$$
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Proof Let $S = T_{F_b}$. For each m in $\mathcal{M} \setminus \mathbb{D}$,

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.

By Hedenmalm's result above, we have that for each m in $\mathcal{M}\backslash\mathbb{D}$, either

- (a) $b \circ \phi_m(z) = \lambda_m$ for some unimodular constant λ_m or
- (b) $\tau_m = b \circ \phi_m(z) \in Aut(\mathbb{D}).$

(A) $b \circ \phi_m(z) = \lambda_m$

 S_m equals the operator $F(\lambda_m)I$ and hence $\sigma(S_m)$ equals one point $F(\lambda_m)$. Thus

$$\sigma(S_m) \subset F(\partial \mathbb{D}) \subset \sigma(T_F),$$

(A) $b \circ \phi_m(z) = \lambda_m$

 S_m equals the operator $F(\lambda_m)I$ and hence $\sigma(S_m)$ equals one point $F(\lambda_m)$. Thus

$$\sigma(S_m) \subset F(\partial \mathbb{D}) \subset \sigma(T_F),$$

and for each λ not in $\sigma(T_F)$,

$$\|(S_m - \lambda I)^{-1}\| = \frac{1}{|F(\lambda_m) - \lambda|}$$

$$\leq \frac{1}{dis(\lambda, \sigma(T_F))}.$$

(B)
$$\tau_m = b \circ \phi_m(z) \in Aut(\mathbb{D})$$

$$S_m = T_{F \circ \tau_m}$$

$$= V_m T_F V_m^*$$

where V_m is the unitary operator on the Bergman space L_a^2 given by

$$V_m f(z) = f(\tau_m(z)) \tau'_m(z).$$

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Thus
$$\sigma(S_m) = \sigma(T_F)$$
 and for each λ in $\mathbb{C} \setminus \sigma(S_m)$,
$$\|(S_m - \lambda I)^{-1}\| = \|V_m T_{F-\lambda}^{-1} V_m^*\|$$
$$= \|T_{F-\lambda}^{-1}\|.$$

So we have

$$\cup_{m\in\mathcal{M}\setminus\mathbb{D}}\sigma(S_m)=\sigma(T_F)$$

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and for each $\lambda \notin \sigma(T_F)$,

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By the following theorem

THEOREM

If $S \in \mathcal{T}$, the Toeplitz algebra, then

$$\mathbb{C}\backslash\sigma_e(S)=\{\lambda\in\mathbb{C}:\lambda
otin \bigcup_{m\in\mathcal{M}\setminus\mathbb{D}}\sigma(S_m) \ \ ext{and}$$

$$\sup_{m\in\mathcal{M}\setminus\mathbb{D}}\|(S_m-\lambda I)^{-1}\|<\infty\}.$$

we have that

$$\sigma_{\mathbf{e}}(T_{F_b}) = \sigma_{\mathbf{e}}(S) = \sigma(T_F).$$

DISCONNECTED ESSENTIAL SPECTRUM

THEOREM

Let h be $\bar{z} + \phi$ such that $\sigma(T_h)$ is disconnected. Let b be an infinite thin Blaschke product and $h_b = h \circ b$. Then

$$\sigma_e(T_{h_b}) = \sigma(T_h)$$

is disconnected.

LEMMA

For each 0 < r < 1, there exists a rational function $\phi(z)$ with poles outside $\overline{\mathbb{D}}$ such that

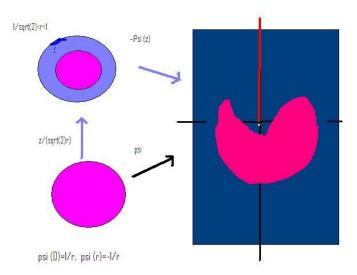
- (a) $2\phi(r) + r\phi'(r) = 0$.
- (b) $1+z\phi(z)$ has a simple zero at z=r and no other zeros in $\overline{\mathbb{D}}$.
- (c) The winding number

$$n(h(\partial \mathbb{D}), 0) = 0$$

where $h = \overline{z} + \phi(z)$.

Proof: For $\frac{1}{\sqrt{2}} < r < 1$, we are going to construct ϕ by some conformal mappings.

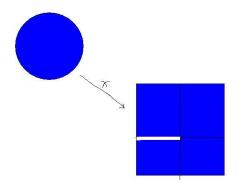
PROOF OF LEMMA



Proof of Lemma

Let λ be the unimodular constant $i\frac{2+i}{2-i}\frac{\sqrt{2}}{1+i}.$ Define

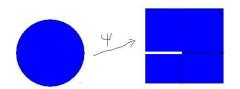
$$\chi(z) = \frac{1}{2r}(\frac{1+z}{1-z})^2,$$



$$\Psi(z) = \chi(\frac{\lambda z - \frac{i}{2-i}}{1 + \frac{i}{2+i}\lambda z}).$$

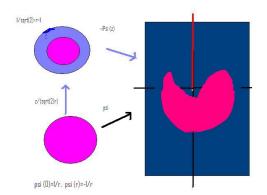
Then

$$\Psi(0) = -\frac{i}{r}, \quad \Psi(\frac{1}{\sqrt{2}}) = \frac{i}{r}.$$



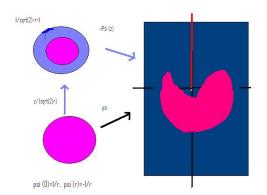
Now define

$$\psi(z)=-i\Psi(\frac{1}{\sqrt{2}r}z).$$



Now define

$$\psi(z) = -i\Psi(\frac{1}{\sqrt{2}r}z).$$



Since $r > \frac{1}{\sqrt{2}}$, the poles of $\psi(z)$ are outside $\overline{\mathbb{D}}$.

Since χ is a conformal map of $\mathbb D$ onto $\mathbb C\setminus (-\infty,0]$, ψ is a conformal map of $\mathbb D$ onto a region bounded by a simple closed curve and 0 is outside the region. In particular $\psi(\partial \mathbb D)$ does not wind around 0 and $\psi(z) \neq 0$ for all z in $\overline{\mathbb D}$.

Since χ is a conformal map of $\mathbb D$ onto $\mathbb C\setminus (-\infty,0], \, \psi$ is a conformal map of $\mathbb D$ onto a region bounded by a simple closed curve and 0 is outside the region. In particular $\psi(\partial \mathbb D)$ does not wind around 0 and $\psi(z) \neq 0$ for all z in $\overline{\mathbb D}$. Defining

$$\phi(z) = \frac{(z-r)\psi(z)-1}{z},$$

we see that (a) and (b) are satisfied:

(a)
$$2\phi(r) + r\phi'(r) = 0$$
.

(b) $1+z\phi(z)$ has a simple zero at z=r and no other zeros in $\overline{\mathbb{D}}$.

On $\partial \mathbb{D}$

$$\overline{z} + \phi(z) = \frac{1}{z} + \phi(z)$$

$$= \frac{1 + z\phi(z)}{z}$$

$$= \frac{z - r}{z}\psi(z).$$

So (c) is satisfied too.

PROOF OF $T_{\bar{z}}f(z) = \frac{1}{z^2} \int_0^z w f'(w) dw$

LEMMA

For f in the Bergman space L_a^2 ,

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Proof of $T_{\bar{z}}f(z) = \frac{1}{z^2} \int_0^z w f'(w) dw$

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For f in the Bergman space L_a^2 ,

$$T_{\bar{z}}f(z)=\frac{1}{z^2}\int_0^z wf'(w)dw.$$

Proof. Note that

$$\{e_n = \sqrt{n+1}z^n\}_{n=0}^{\infty}$$

is an orthonormal basis of the Bergman space. To prove this lemma, we need only verify the above equality for each $f(z) = e_n$.

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is an orthonormal basis of the Bergman space. To prove this lemma, we need only verify the above equality for each $f(z)=e_n$. As $T_{\bar{z}}$ is the adjoint of the Bergman shift, we have

$$T_{\overline{z}}e_n = \begin{cases} 0 & n = 0\\ \sqrt{\frac{n}{n+1}}e_{n-1}, & n > 0 \end{cases}$$

On the other hand, since $e_n(w) = \sqrt{n+1}w^n$, an easy calculation gives

$$\int_0^z we_n'(w)dw = \frac{nz^{n+1}}{\sqrt{n+1}}$$

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Thus we have

$$\frac{1}{z^{2}} \int_{0}^{z} we'_{n}(w) dw = \frac{nz^{n-1}}{\sqrt{n+1}}$$
$$= \sqrt{\frac{n}{n+1}} e_{n-1},$$

to obtain

$$T_{\bar{z}}e_n=\frac{1}{z^2}\int_0^z we'_n(w)dw.$$

This completes the proof of the lemma.